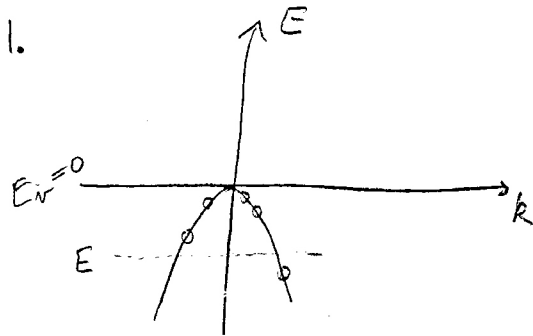


ECE G287 Optical Detectors HW2 Solution



$$p = \int_{-\infty}^{E_v} (1 - f(E)) N(E) dE$$

Band energy with $E_v = 0$

$$E = -\frac{\hbar^2 k^2}{2m^*}$$

$$N_T(E) = \# \text{ of valence band states with } E' \geq E$$

$$= \frac{2}{(2\pi)^3} \overset{\substack{\text{spin} \\ \nearrow}}{4\pi} \overset{\substack{\text{K-space volume} \\ \nearrow \text{ of one state}}}{K(E)}^3$$

where $K(E) = \sqrt{-\frac{2m^*E}{\hbar^2}}$

$$\Rightarrow N_T(E) = \frac{2}{(2\pi)^3} \frac{4\pi}{3} \left(-\frac{2m^*E}{\hbar^2} \right)^{3/2}$$

$$N(E) = \text{Density of states} = \frac{dN_T(E)}{dE} = \frac{2}{(2\pi)^3} \frac{4\pi}{3} \left(\frac{2m^*}{\hbar^2} \right)^{3/2} \frac{3}{2} (-E)^{1/2}$$

$$= \frac{2}{(2\pi)^2} \left(\frac{2m^*}{(\hbar/2\pi)^2} \right)^{3/2} (-E)^{1/2}$$

$$1 - f(E) = \frac{1}{e^{(E-E_f)/kT} + 1}$$

$$= \frac{e^{-(E-E_f)/kT}}{e^{(E-E_f)/kT} + 1} = \frac{1}{e^{-(E-E_f)/kT} + 1}$$

Since for the entire valence band $E < 0$ and $-(E-E_f) \gg kT$ can neglect the 1 in the denominator.

$$\Rightarrow \rho \approx \int_{-\infty}^{\infty} \frac{1}{e^{-(E-E_f)/kT}} \frac{2}{(2\pi)^2} \left(\frac{2m^*}{(\hbar/2\pi)^2} \right)^{3/2} (-E)^{1/2} dE$$

Changing integration variable from $E \rightarrow -E$ and reversing limits of integration:

$$\rho = \int_0^{\infty} e^{-(E+E_f)/kT} \frac{2}{(2\pi)^2} \left(\frac{2m^*}{(\hbar/2\pi)^2} \right)^{3/2} E^{1/2} dE$$

$$= e^{-E_f/kT} \frac{2}{(2\pi)^2} \left(\frac{2m^*}{(\hbar/2\pi)^2} \right)^{3/2} \int_0^{\infty} e^{-E/kT} E^{1/2} dE$$

Definite integral is given by: $\int_0^{\infty} x^{1/2} e^{-ax} dx = \frac{\sqrt{\pi}}{2a\sqrt{a}}$

$$\Rightarrow \rho = e^{-E_f/kT} \frac{2}{(2\pi)^2} \left(\frac{2m^*}{(\hbar/2\pi)^2} \right)^{3/2} \frac{\sqrt{\pi} (kT)^{3/2}}{2}$$

$$= \frac{1}{\sqrt{2}} \frac{1}{(2\pi)^{3/2}} \left(\frac{2m^* kT}{(\hbar/2\pi)^2} \right)^{3/2} e^{-E_f/kT}$$

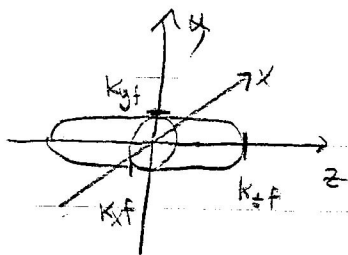
$$= 2 \left(\frac{2\pi m^* kT}{h^2} \right)^{3/2} e^{-E_f/kT}$$

If we change the energy scale such that $E_v \neq 0$, have to express all energies as $E - E_v$!

$$\rho = 2 \left(\frac{2\pi m^* kT}{h^2} \right)^{3/2} e^{-(E_f - E_v)/kT}$$

$$2. \quad E = \frac{(\hbar k_z)^2}{2m_l} + \frac{(\hbar k_x)^2}{2m_t} + \frac{(\hbar k_y)^2}{2m_t}$$

Constant energy surface is ellipsoid in k -space



$$k_{yf} = k_{xf} = \sqrt{\frac{2m_t E}{\hbar^2}}$$

$$k_{zf} = \sqrt{\frac{2m_l E}{\hbar^2}}$$

$$\text{Volume} = \frac{4\pi}{3} k_{xf} k_{yf} k_{zf} = \frac{4\pi}{3} \left(\frac{2E}{\hbar^2}\right)^{3/2} (m_t^2 m_l)^{1/2}$$

$$N_T(E) = \# \text{ of states with } E' < E$$

$$= 2 \cdot 6 \cdot \frac{1}{(2\pi)^3} \frac{4\pi}{3} \left(\frac{2E}{\hbar^2}\right)^{3/2} (m_t^2 m_l)^{1/2}$$

Spin states

number of conduction band pockets

k -space volume per electron state

$$N(E) = \frac{dN_T(E)}{dE} = \frac{2}{(2\pi)^3} \frac{4\pi}{3} \left[6 (m_t^2 m_l)^{1/2} \right] \left(\frac{2}{\hbar^2}\right)^{3/2} \frac{3}{2} E^{1/2}$$

Comparing with Problem 1, we see that:

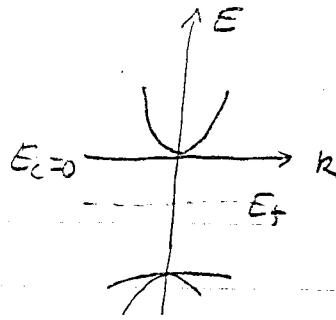
$$6 (m_t^2 m_l)^{1/2} = (m_d^*)^{3/2}$$

$$\Rightarrow m_d = 6^{2/3} (0.98 \cdot 0.19 \cdot 0.19)^{1/3} m_0$$

$$= 1.08 m_0$$

The rest of the derivation follows as in Problem 1

Setting $E_c = 0$



$$n = \int_0^{\infty} \frac{1}{e^{(E-E_f)/kT} + 1} \cdot \frac{2}{(2\pi)^2} \left(\frac{2m_d^*}{\hbar^2} \right)^{3/2} E^{1/2} dE$$

$$\text{For } E - E_f \gg kT: \frac{1}{e^{(E-E_f)/kT} + 1} \approx \frac{1}{e^{(E-E_f)/kT}} = e^{-(E-E_f)/kT}$$

$$n = e^{E_f/kT} \cdot \frac{2}{(2\pi)^2} \left(\frac{2m_d^*}{\hbar^2} \right)^{3/2} \int_0^{\infty} E^{1/2} e^{-E/kT} dE$$

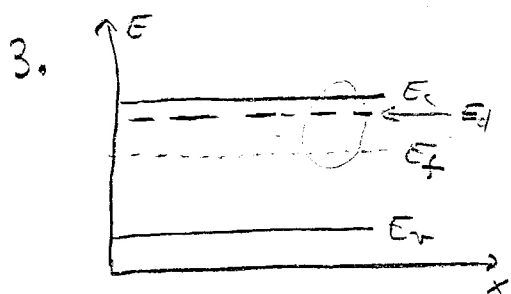
$$= e^{E_f/kT} \frac{2}{(2\pi)^2} \left(\frac{2m_d^*}{\hbar^2} \right)^{3/2} \frac{\sqrt{\pi} (kT)^{3/2}}{2}$$

$$\Rightarrow n = 2 \left[\frac{2\pi m_d^* kT}{\hbar^2} \right]^{3/2} e^{E_f/kT}$$

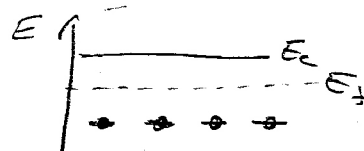
If we let $E_c \neq 0$, we have to set $E \rightarrow E - E_c$

$$\Rightarrow n = 2 \left[\frac{2\pi m_d^* kT}{\hbar^2} \right]^{3/2} e^{-(E_c - E_f)/kT}$$

$$\text{where } m_d^* = 1.08 m_0$$



At $T=0$, Fermi level is between donor level and conduction band edge:



Set $E_v = 0$

$$\Rightarrow E_c = E_g = 1.11 \text{ eV}$$

$$E_d = E_c - 0.044 \text{ eV} = 1.066 \text{ eV}$$

Charge balance

$\Rightarrow$ # of conduction-band electrons

= # of unoccupied donor sites + # of holes

$$n = N_d^+ + p$$

$$= N_d (1 - f(E_d)) + p$$

$$n = 2 \left(\frac{2\pi m_n^* kT}{h^2} \right)^{3/2} e^{-(E_c - E_f)/kT}$$

$$= (1.1)^{3/2} (T)^{3/2} (4.818 \times 10^{21}) e^{-(E_c - E_f)/kT}$$

$$p = 2 \left(\frac{2\pi m_p^* kT}{h^2} \right)^{3/2} e^{-(E_f - E_v)/kT}$$

$$= (0.55)^{3/2} (T)^{3/2} (4.818 \times 10^{21}) e^{-(E_f - E_v)/kT}$$

$$N_d^+ = N_d \left(1 - \frac{1}{e^{(E_d - E_f)/kT} + 1} \right) = N_d \left[\frac{1}{e^{-(E_d - E_f)/kT} + 1} \right]$$

$$= 3 \times 10^{23} \frac{1}{e^{-(E_d - E_f)/kT} + 1}$$

Charge balance $\Rightarrow$

$$\begin{aligned} (1.1)^{3/2} (T)^{3/2} &= (4.818 \times 10^{21}) e^{-(E_c - E_f)/kT} \\ &= (0.55)^{3/2} (T)^{3/2} (4.818 \times 10^{21}) e^{-(E_f - E_v)/kT} \\ &\quad + 3 \times 10^{23} [e^{-(E_c - E_f)/kT} + 1]^{-1} \end{aligned}$$

Divide by 10^{21} and put in values for E_c , E_d , and E_v ($E_v = 0$)

$$\begin{aligned} (4.818) (T^{3/2}) &[(1.1)^{3/2} e^{-(1.11 - E_f)/kT} - (0.55)^{3/2} e^{-E_f/kT}] \\ &- 300 [e^{-(1.066 - E_f)/kT} + 1]^{-1} = 0 \end{aligned}$$

Define the left-hand side of equation above as $f(E_f)$.
Numerically solve to see what value of E_f makes $f(E_f) = 0$.

C:\MATLAB6p5\work\fermi_level.m

March 16, 2004

```
function y=fermi_level(Ef)
T=4.2;
kT=T*1.38e-23/1.6e-19;
exp1=exp( -(1.11-Ef)/kT );
exp2=exp( -Ef/kT );
exp3=exp( -(1.066-Ef)/kT );
y=4.818*T^1.5 *(1.1^1.5 *exp1 - 0.55^1.5 *exp2) - 300 ./ (exp3 +1);
```

a) $T = 4.2\text{K} \Rightarrow kT = 0.000362\text{ eV}$

From MATLAB, $E_f = 1.0833$

$E_c - E_f = 0.0217\text{ eV}$ $(E_c - E_f)/kT = 59.8$

b) $T = 77\text{K} \Rightarrow kT = 0.00664\text{ eV}$

From MATLAB: $E_f = 1.0799$; $E_c - E_f = 0.0308 = 4.64\text{ kT}$

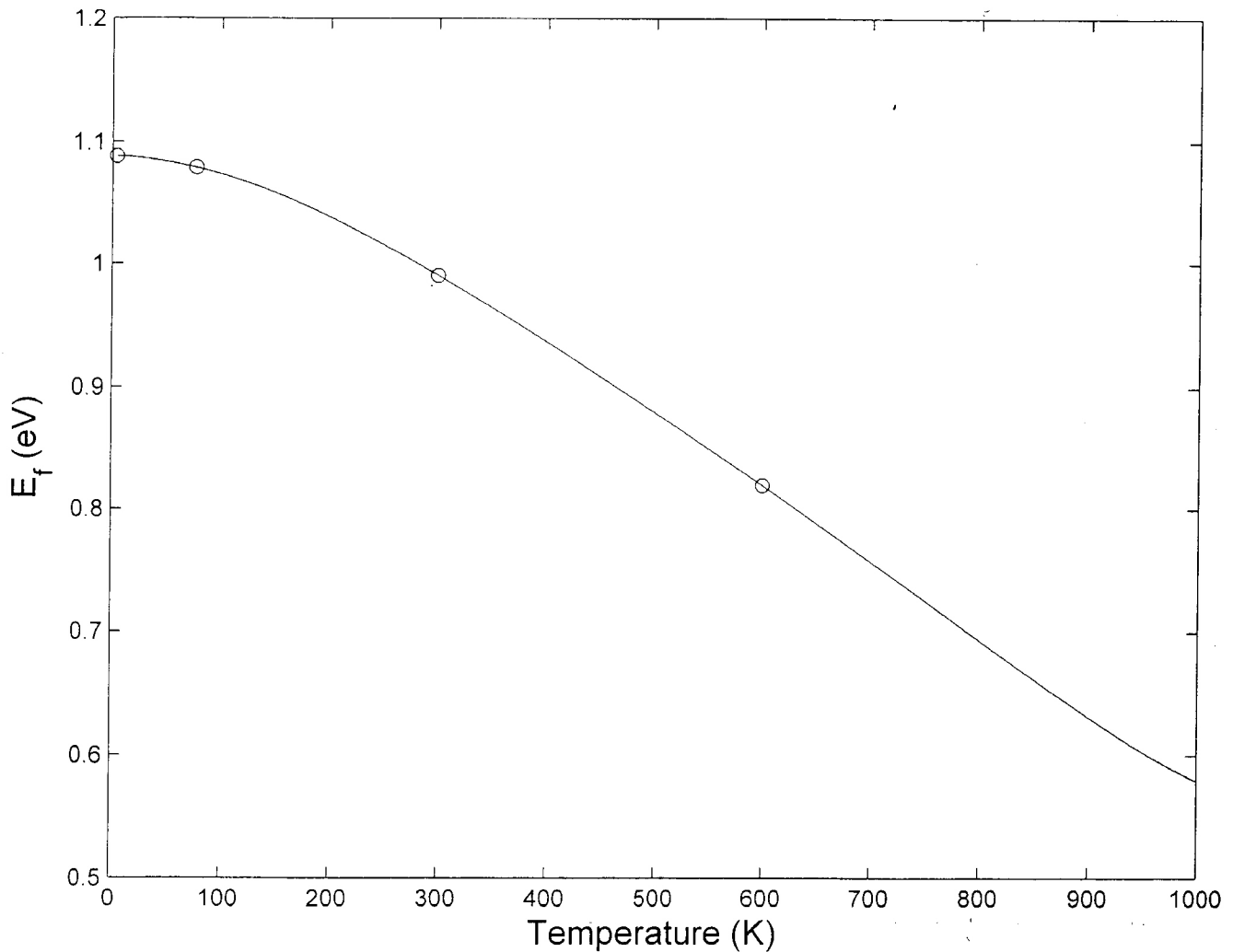
c) $T = 300\text{K} \Rightarrow kT = 0.0259\text{ eV}$

From MATLAB: $E_f = 0.9905$ $(E_c - E_f)/kT = 9.62$

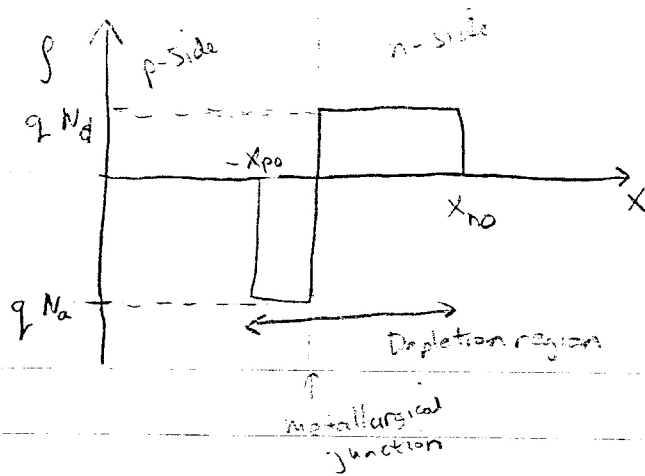
d) $T = 600\text{K} \Rightarrow kT = 0.0517$

From MATLAB: $E_f = 0.8194$ $(E_c - E_f)/kT = 5.62$

```
>> global T
>> T=4.2
T =
    4.2000
>> Ef42=fzero('fermi_level', 1.0)
Ef42 =
    1.0883
>> T=77
T =
    77
>> Ef77=fzero('fermi_level', 1.0)
Ef77 =
    1.0792
>> T=300
T =
    300
>> Ef300=fzero('fermi_level', 1.0)
Ef300 =
    0.9905
>> T=600
T =
    600
>> Ef600=fzero('fermi_level', 1.0)
Ef600 =
    0.8194
>>
```



4. Abrupt junction:

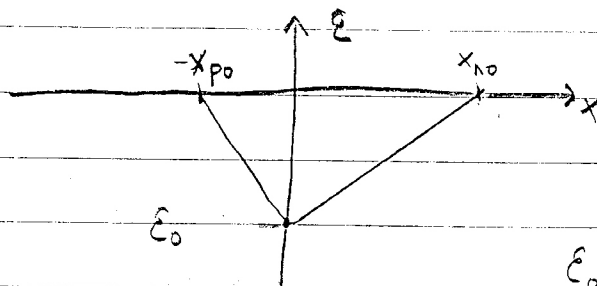


$$\text{Stored charge} = Q_0 = AqN_ax_{op} = AqN_dx_{on}$$

Relation of junction width to applied voltage:

Poisson's equation: $\nabla \cdot \epsilon \mathcal{E} = \rho$
 $\epsilon \frac{d\mathcal{E}}{dx} = \rho$

Integrating across junction:



$$\mathcal{E}_0 = -\frac{q}{\epsilon} N_a x_{po} = -\frac{q}{\epsilon} N_d x_{no}$$

For $x < 0$:

$$\mathcal{E}(x) = \int_{-x_{po}}^x -\frac{q}{\epsilon} N_a dx = -\frac{q}{\epsilon} N_a (x + x_{po})$$

For $x > 0$:

$$\mathcal{E}(x) = \int_{x_{no}}^x \frac{q}{\epsilon} N_d dx = \frac{q}{\epsilon} N_d (x - x_{no})$$

$$E(x) = - \frac{dV(x)}{dx}$$

$$\Rightarrow \Delta V = - \int E(x) dx$$

Integrating across entire junction yields:

$$\Delta V = -\frac{1}{2} E_0 (x_{p0} + x_{n0})$$

$$= \frac{q N_a}{2 \epsilon} x_{p0} (x_{p0} + x_{n0})$$

Under bias:

$$\Delta V \rightarrow V_0 - V_a$$

$$x_{p0} \rightarrow x_p$$

$$x_{n0} \rightarrow x_n$$

$$V_0 - V_a = \frac{q}{2} \frac{N_a}{\epsilon} x_p (x_p + x_n)$$

$$\text{Stored charge} = Q = A q N_a x_p = A q N_d x_n$$

$$\Rightarrow x_n = \frac{N_a}{N_d} x_p$$

$$V_0 - V_a = \frac{q}{2} \frac{N_a}{\epsilon} x_p \left(x_p + \frac{N_a}{N_d} x_p \right) = \frac{q}{2} \frac{N_a}{\epsilon} x_p^2 \left(1 + \frac{N_a}{N_d} \right)$$

$$2 \epsilon (V_0 - V_a) \left(1 + \frac{N_a}{N_d} \right)^{-1} = \frac{1}{q N_a A^2} Q^2$$

$$Q = \sqrt{2 \epsilon q N_a A^2 (V_0 - V_a) \left(1 + \frac{N_a}{N_d} \right)^{-1}}$$

$$C_j = \frac{dQ}{dV_a} = \frac{dQ}{d(V_0 - V_a)} \frac{d(V_0 - V_a)}{dV_a} = - \frac{dQ}{d(V_0 - V_a)} = \frac{1}{2} \frac{1}{Q} \left(2 \epsilon q N_a A^2 \left(1 + \frac{N_a}{N_d} \right)^{-1} \right)$$

$$C_j = \frac{1}{2} \frac{2 \epsilon q N_a A^2 \left(1 + \frac{N_a}{N_d} \right)^{-1}}{\sqrt{2 \epsilon q N_a A^2 (V_0 - V_a) \left(1 + \frac{N_a}{N_d} \right)^{-1}}} = \frac{A}{2} \sqrt{\frac{2 q \epsilon N_a \left(1 + \frac{N_a}{N_d} \right)^{-1}}{(V_0 - V_a)}}$$

$$C_j = \frac{A}{2} \sqrt{\frac{2 q \epsilon}{(V_0 - V_a)} \frac{N_a N_d}{(N_a + N_d)}}$$

Approximate parallel-plate capacitor solution:

$$C' = \epsilon \frac{A}{d} = \epsilon \frac{A}{W}$$

$$\Delta V = (V_0 - V_a) = \frac{q N_a}{2\epsilon} x_p (x_p + x_n)$$

$$(x_p + x_n) = W$$

$$x_p \left(1 + \frac{N_a}{N_d}\right) = W \Rightarrow x_p = W \left(1 + \frac{N_a}{N_d}\right)^{-1}$$

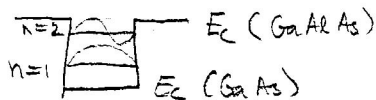
$$\Rightarrow (V_0 - V_a) = \frac{q N_a}{2\epsilon} W^2 \left(1 + \frac{N_a}{N_d}\right)^{-1}$$

$$W = \left[\frac{2\epsilon (V_0 - V_a)}{q \frac{N_a N_d}{(N_a + N_d)}} \right]^{1/2}$$

$$C' = \frac{\epsilon A}{\sqrt{\frac{2\epsilon (V_0 - V_a)}{q \frac{N_a N_d}{(N_a + N_d)}}}} = A \sqrt{\frac{\epsilon q}{2(V_0 - V_a)} \frac{N_a N_d}{(N_a + N_d)}}$$

This is the same as the exact result!

S.



$$E = E_c + E_n + \frac{(\hbar k_x)^2}{2m^*} + \frac{(\hbar k_y)^2}{2m^*}$$

$$\text{where } E_n = \frac{\hbar^2}{2m^*} (n\pi/d)^2 \quad n=1, 2, 3, \dots$$

Setting $E_c = 0$:

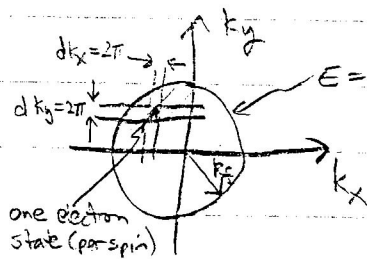
$$E = E_n + \frac{\hbar^2}{2m^*} (k_x^2 + k_y^2) = \frac{\hbar^2}{2m^*} \left[\left(\frac{n\pi}{d} \right)^2 + k_t^2 \right]$$

(See plot)

Density of states ($n=1$)

$$N_1(E) = 0 \text{ for } E < E_1 = \frac{\hbar^2}{2m^*} \left(\frac{\pi}{d} \right)^2$$

For $E > E_1$, need to count states in k -space



$$E = E_1 = \frac{\hbar^2}{2m^*} (k_x^2 + k_y^2) + \frac{\hbar^2}{2m^*} \left(\frac{n\pi}{d} \right)^2$$

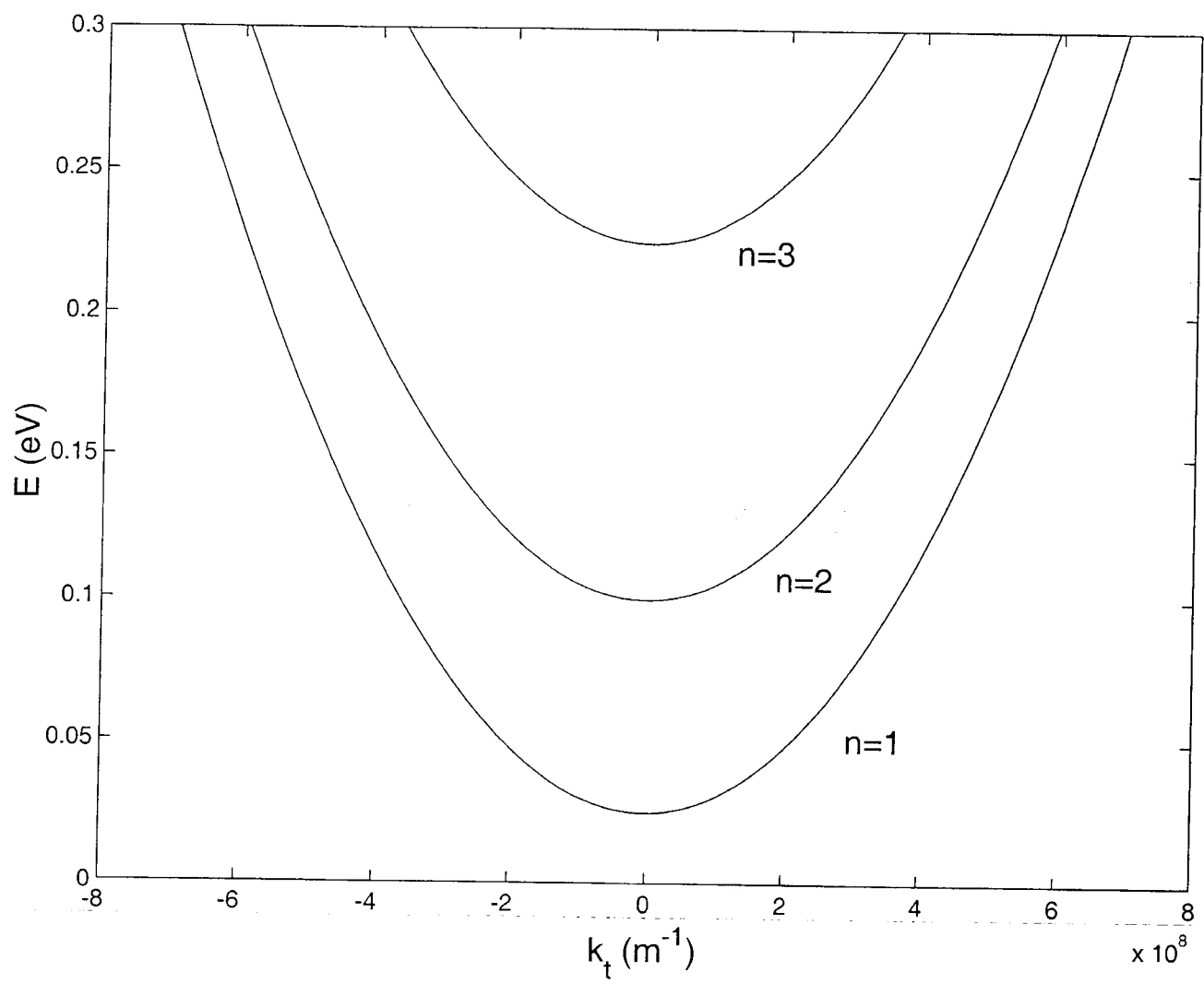
$$k_f = \sqrt{\frac{2m^*}{\hbar^2} (E - E_n)}$$

N_T = total number of states with $E < E'$

$$= \frac{2}{(2\pi)^2} \pi (k_f)^2 = \frac{2\pi}{(2\pi)^2} \left[\frac{2m^*}{\hbar^2} E - \left(\frac{n\pi}{d} \right)^2 \right]$$

2 Spin states

$$N_1(E) = \frac{dN_T}{dE} = \frac{2\pi}{(2\pi)^2} \frac{2m^*}{\hbar^2} \text{ for } E > E_1$$



Density of state ($n=2$)

$$N_2(E) = 0 \text{ for } E < E_2 \quad (E_2 = \frac{\hbar^2}{2m^*} \left(\frac{2\pi}{d}\right)^2)$$

$$N_T = \frac{2}{(2\pi)^2} \pi (k_{f2})^2 = \frac{2\pi}{(2\pi)^2} \left[\frac{2m^*}{\hbar^2} E - \left(\frac{2\pi}{d}\right)^2 \right]$$

$$N_2(E) = \frac{dN_T}{dE} = \frac{2\pi}{(2\pi)^2} \frac{2m^*}{\hbar^2} \text{ for } E > E_2$$

$$\text{Similarly: } N_3(E) = \frac{2\pi}{(2\pi)^2} \frac{2m^*}{\hbar^2} \text{ for } E > E_3$$

$$N(E) = N_1 + N_2 + N_3 + \dots$$

$$\text{Units: } N(E) \leftrightarrow \frac{(\text{Kg})}{(\text{J} \cdot \text{sec})^2} = \frac{(\text{Kg})}{\text{J}^2 \text{ sec}^2} = \frac{(\text{Kg})}{\text{Kg m}^2/\text{sec}^2 \cdot \text{sec}^2 \text{ J}} = \text{m}^{-2} \text{J}^{-1}$$

$$= \text{electron states/m}^2/\text{J}$$

which is correct for 2D density of states

(To convert to electron states/m²/eV, multiply by 1.6×10^{-19})

